

hemisphere, as seen in Fig. 2c. This streaking precipitation is produced by a combination of convergent and circulating surface winds, indicating a role for both gravity waves and Rossby waves. The depth of convection in the streaking feature is shallower than the chevron (Supplementary Fig. S6), however its lifetime is considerably longer, resulting in several centimetres of precipitation over its extent. Currently published ISS images do not show evidence for surface changes associated with the storm in Fig. 1f, which suggests these clouds do not produce substantial rainfall. This indicates that our model overestimates the amount of mid-latitude precipitation during the current season.

Precipitation rates within these modes exceed the zonal- and time-mean rate by a factor of ~ 20 during the equinoctial transition (Fig. 3a). We infer that three-dimensional dynamics are responsible for local and significant enhancements in precipitation, and are therefore essential for understanding the observed distribution of surface erosion²⁹. Equatorial chevrons are ephemeral features of Titan's climate, however, occurring on the order of an Earth year as Titan's ITCZ undergoes a post-equinoctial transition from one hemisphere to the other (Supplementary Figs S7 and S10; ref. 17). Surface accumulation (that does not infiltrate into the regolith) evaporates away during the summer (Supplementary Fig. S8), with important implications for observed fluvial features in Titan's semi-arid environment²⁹.

We now show that the cloud structures in the examples above are not isolated occurrences, but instead constitute the 'typical' behaviour of convection in the model, organized by well-defined dynamical modes of the atmosphere. The evolution in time and longitude of modelled equatorial precipitation (Fig. 3a) shows evidence of eastward-propagating disturbances with a phase speed of about 12 m s^{-1} (indicated by the solid black line), superposed on slower disturbances travelling westwards at around 1 m s^{-1} (dashed line). These two modes of variability are also apparent in a space-time power spectrum of equatorial surface zonal wind (Fig. 3b), which shows sharp peaks at an eastward-propagating zonal wavenumber of 2 with a period of about eight days (corresponding to a phase speed of about 12 m s^{-1} at the equator) and at a westward-propagating zonal wavenumber of 2 with period ~ 100 days (phase speed $\sim 1 \text{ m s}^{-1}$). This close correspondence between precipitation and surface wind perturbations implies a link between dynamical and convective processes, as seen in convectively coupled equatorial waves on Earth¹⁹.

To isolate the spatial structure of the dominant eastward and westward modes, we follow a three-step procedure described in Methods. The spatial structure and phase speed of the eastward-propagating mode corresponds closely to that of a baroclinic, equatorially trapped Kelvin wave (see Supplementary Information). The spatial patterns of precipitation and surface wind anomaly of this mode are shown in Fig. 2c. Crucially, these waves are associated with roughly chevron-shaped equatorial precipitation patterns. The westward-propagating mode (Fig. 2d) stretches from the equator to the high-latitudes of the Southern Hemisphere, and is associated with a streak-like precipitation feature stretching southeastwards from the tropics into the mid-latitudes. Thus, convective storms in the GCM are organized by large-scale modes of variability, and the modes are robustly associated in a statistical sense with precipitation structures resembling the observed clouds.

We have also conducted sensitivity tests to investigate the feedback of moist convective heating onto the dynamical modes (see Supplementary Information for details). In a simulation where latent heating by convection and surface evaporation is artificially suppressed, we find that the eastward-propagating Kelvin mode is unchanged, whereas the westward mode is substantially altered (Fig. 3c). Thus the Kelvin mode is convectively uncoupled (see Supplementary Information), although its surface convergence field plays an important role in organizing precipitation. Convection

Figure 3 | Space-time variability of simulated equatorial winds and precipitation averaged over $\pm 10^\circ$ latitude in the Titan GCM.

a, Longitude-time distribution of precipitation rate normalized to the global- and time-mean rate for the period containing the chevron event in Fig. 1 (magenta triangle at day 30) and the streak event (triangle at day 210). Lines indicate constant velocity trajectories of 11.7 m s^{-1} eastwards (solid) and 0.9 m s^{-1} westwards (dashed). **b**, Space-time spectral decomposition of surface zonal winds for 1,000 simulation days centred around the equinox. **c**, As in **b** but for a test simulation with the latent heating of methane removed.

is essential to the dynamics of the westward mode, however, as indicated by the substantial reduction of precipitation at negative wavenumbers in the test case (Fig. 3c; also see Supplementary Information). This is the first evidence of a convectively coupled wave on a planetary body other than Earth.

A sequence of cloud observations taken in 2008 (before NSE) with ground-based instruments indicated a connection between disturbances at equatorial and polar latitudes likely to be mediated by planetary-scale waves¹¹. This cloud propagated eastwards at $\sim 3 \text{ m s}^{-1}$, and was interpreted as a stationary Rossby wave advected by the mid-tropospheric wind. However, our simulation indicates that the eastward-propagating precipitation field during this epoch is associated with the Kelvin wave component. As this observation occurred before NSE, surface convergence resulting

from the superposition of the Kelvin mode and the ITCZ is most intense south of the equator (Supplementary Fig. S10). As the season transitions towards Northern Summer Solstice (NSS), our model indicates a shift in the precipitation by the Kelvin–ITCZ superposition to northern mid-latitudes accompanied by southern mid-latitude streaks. We therefore predict a phase of transient cloud activity in both hemispheres of Titan during the next several years (Supplementary Fig. S7).

In summary, we have developed a process for interpreting the morphologies of and precipitation associated with Titan's clouds through a combined analysis of observations and GCM simulations, thereby opening a new field of dynamic meteorology. We find that recent cloud activity near Titan's equator just following NSE is consistent with the presence of two dominant modes of atmospheric variability in the GCM. A fast, eastward-propagating mode travelling at $\sim 12 \text{ m s}^{-1}$ with the character of an equatorially trapped Kelvin wave produces chevron-shaped precipitation patterns similar to the clouds observed by Cassini ISS on 27 September 2010 (ref. 9). A slow, westward-propagating mode that is vertically confined accounts for streaks of precipitation such as the cloud observed on 18 October 2010 (ref. 9). This latter mode is convectively coupled, the first of its kind to be inferred outside of Earth's atmosphere. The phasing of these modes produce several centimetres of precipitation over 1,000-km-scale regions, locally enhancing precipitation rates by 20-fold over the mean. These modes therefore play a crucial role in fluvial erosion of Titan's surface. Observations clearly indicate surface changes associated with the equatorial chevron¹⁰, but precipitation from the mid-latitude streak is too light to be detected, indicating the model overestimates mid-latitude precipitation during the current season.

Methods

Our Titan GCM is similar to that used in a previous study¹⁶, the primary difference being that the model is now fully three-dimensional. Simplified treatments of radiation, convection, and the seasonal cycle are included. The model does not include tides, either thermal (diurnal) or gravitational (semi-diurnal), nor any other non-axisymmetric forcing mechanism (be it topography, albedo, opacity sources, or others.). The atmosphere is spun up and moistened by methane evaporation from an initial state of solid-body rotation and zero methane humidity. The seasonal cycle of insolation is included in the model forcing; see the Supplementary Information for a more detailed description. A (statistical) steady state is achieved after ~ 20 Titan years (~ 600 Earth years). The results shown in the figures are from 1,000 terrestrial days bracketing NSE of the 21st simulated Titan year. (The zonal- and time-mean overturning streamfunction, zonal winds, and temperature of the full, 1,800 model days are shown in the Supplementary Information.)

The GCM does not include a cloud scheme, so we have developed a method for simulating cloud observations based on precipitation. Briefly, we assume cloud droplets are distributed about a peak size of $512 \mu\text{m}$, which are marginally lofted for convective updrafts of 10 m s^{-1} , typical for deep convective events on Titan²¹. We assume 10–20% of the column precipitation remains suspended, with a cloud-top altitude of 20 km. Together the column mass of the cloud and the droplet size distribution determine the cloud optical depth. Mie theory then translates the droplet size distribution into the scattering properties of clouds, ultimately linking the precipitation field from the GCM to the radiative transfer model (see Supplementary Information for more details).

To isolate the spatial structure of dominant modes of variability in the GCM, the surface zonal wind field is spectrally filtered to retain only the space and timescales of the dominant spectral peaks; the leading patterns of variability of the filtered surface wind are extracted using empirical orthogonal function (EOF) analysis, and the associated three-dimensional fields are reconstructed by regression onto the leading principal component; further details are given in the Supplementary Information.

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Author contributions

J.L.M. contributed experiment design, performed GCM simulation analysis, and wrote the manuscript. M.Á. and J.L.M. developed the algorithm for simulating clouds from the GCM, and M.Á. contributed radiative transfer analysis of the GCM. R.C. contributed the statistical analysis of the GCM simulation and J.L.M. and R.C. interpreted the model dynamics. E.P.T. contributed analysis of Cassini ISS cloud images.

Additional information

The authors declare no competing financial interests. Supplementary information accompanies this paper on www.nature.com/naturegeoscience. Reprints and permissions information is available online at <http://www.nature.com/reprints>. Correspondence and requests for materials should be addressed to J.L.M.