

## A THEOREM ON INFINITE PRODUCTS OF EIGENVALUES OF STURM-LIOUVILLE TYPE OPERATORS

S. LEVIT AND U. SMILANSKY

**ABSTRACT.** Infinite products of ratios of eigenvalues of Sturm-Liouville operators are expressed in a closed form in terms of corresponding solutions of initial-value problems.

**Introduction.** Gaussian path integrals are defined by

$$(1) \quad I = \int_{\text{paths}} D[\vec{\eta}(t)] \exp[iS[\vec{\eta}(t)]]$$

where

$$(2) \quad S[\vec{\eta}(t)] = \int_0^T (\vec{\eta}(t) \cdot \Lambda(t)\vec{\eta}(t)) dt,$$

and  $\vec{\eta}(t)$  are the  $N$  dimensional paths satisfying the end conditions  $\vec{\eta}(0) = \vec{\eta}(T) = 0$ .  $\Lambda(t)$  is a differential operator of second order in  $(d/dt)$  acting on the space of paths. Such path integrals are often encountered in physical problems, e.g. the formulation of quantum mechanics [1].

The integral (1) can be evaluated in terms of the infinite product of the eigenvalues of  $\Lambda(t)$  [2], [3]. In this note we present a theorem which enables one to express these infinite products in a closed form.

The theorem will be presented with a detailed proof for the simple case of a one dimensional problem ( $N = 1$ ). The extension to  $N > 1$  will be formulated but a detailed proof will not be supplied since it follows the same lines as for  $N = 1$ .

**THEOREM 1.** *Let the differential operator*

$$(3) \quad \Lambda(\alpha; t) = (d/dt)(p(t)d/dt) + \alpha q(t)$$

*be defined for  $0 \leq t \leq T$ ,  $0 \leq \alpha \leq 1$ , where  $p(t)$  and  $q(t)$  are smooth functions of  $t$ ,  $p(t) > p_0 > 0$ .*

*Let  $\lambda_k(\alpha)$  and  $U^{(k)}(\alpha; t)$  be the eigenvalues and eigenfunctions for the boundary-value problem*

$$(4) \quad \Lambda(\alpha; t)U(\alpha; t) + \lambda(\alpha)U(\alpha; t) = 0, \quad U(\alpha; 0) = U(\alpha; T) = 0.$$

*Let  $y(\alpha; t)$  be the solution of the initial-value problem*

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$$(5) \quad \Lambda(\alpha; t)y(\alpha; t) = 0, \quad y(\alpha; 0) = 0, \quad (d/dt)y(\alpha; 0) = 1.$$

Then, for every  $\alpha$

$$(6) \quad \left| \prod_{k=1}^{\infty} \frac{\lambda_k(\alpha)}{\lambda_k(0)} \right| = \left| \frac{y(\alpha; T)}{y(0; T)} \right|.$$

PROOF. The following points should be observed:

(a) For all  $k$ ,  $\lambda_k(0) > 0$  and for large  $k$ ,

$$\lambda_k(0) \cong \left[ \pi / \int_0^T \frac{dt}{\sqrt{p(t)}} \right]^2 \cdot k^2, \quad [4].$$

(b) The functions  $\lambda_k(\alpha)$ ,  $U^{(k)}(\alpha; t)$  and  $y(\alpha; t)$  are analytic in  $\alpha$  [5].

(c) For large  $k$ ,  $|\lambda_k(\alpha) - \lambda_k(0)| \leq \alpha \cdot M$  [4] with  $M$  independent of  $\alpha$  or  $k$ . Hence the product on the left-hand side of (6) converges uniformly for all finite values of  $\alpha$ .

(d) There is a finite number of nonpositive eigenvalues for a given  $\alpha$  [6]. Hence the product on the left-hand side of (6) vanishes only at a finite number of  $\alpha$  values, in the interval  $0 \leq \alpha \leq 1$ .

A necessary and sufficient condition for  $y(\alpha; T)$  to vanish is that one of the  $\lambda_k(\alpha)$  vanishes. Hence the two sides of equation (6) vanish simultaneously.

Consider now the functions:

$$(7) \quad f(\alpha) = \prod_{k=1}^{\infty} \frac{\lambda_k(\alpha)}{\lambda_k(0)},$$

$$(8) \quad \tilde{f}(\alpha) = y(\alpha; T)/y(0; T).$$

Since by construction  $f(0) = \tilde{f}(0) = 1$ , the theorem will be proved once it is demonstrated that  $(d/d\alpha)f(\alpha)/f(\alpha) = (d/d\alpha)\tilde{f}(\alpha)/\tilde{f}(\alpha)$ , for all  $\alpha$  satisfying  $f(\alpha) \neq 0$ .

By virtue of (c) one can differentiate the left-hand side of equation (7) as if it were a finite product and get

$$(9) \quad \frac{d}{d\alpha} f(\alpha)/f(\alpha) = \sum_{k=1}^{\infty} \frac{d\lambda_k(\alpha)}{d\alpha} \frac{1}{\lambda_k(\alpha)}.$$

It can easily be shown that

$$(10) \quad \frac{d\lambda_k(\alpha)}{d\alpha} = - \int_0^T [U^{(k)}(\alpha; t)]^2 q(t) dt,$$

where  $U^{(k)}(\alpha; t)$  are the solutions of (4) subject to the normalization

$$(11) \quad \int_0^T [U^{(k)}(\alpha; t)]^2 dt = 1.$$

Since the integral in (10) is bounded by  $\text{Max}|q(t)|$  and because of (c), the series on the right-hand side of (9) converges uniformly and the summation and integration operations could be interchanged to yield

$$(12) \quad \frac{d}{d\alpha} f(\alpha)/f(\alpha) = \int_0^T G_{\lambda=0}(t, t, \alpha)q(t) dt$$

with

$$(13) \quad G_{\lambda}(t, t', \alpha) = \sum_{k=1}^{\infty} \frac{U^{(k)}(\alpha; t)U^{(k)}(\alpha; t')}{\lambda - \lambda_k(\alpha)},$$

which is the spectral representation of Green's function for the boundary-value problem (4).

Returning to the initial value problem (5) and differentiating it with respect to  $\alpha$ , one gets that  $z(\alpha; t) \equiv dy(\alpha; t)/d\alpha$  is a solution of the initial-value problem

$$(14) \quad \Lambda(\alpha; t)z(\alpha; t) = -q(t)y(\alpha; t), \quad z(\alpha; 0) = 0, \quad (d/dt)z(\alpha; 0) = 0.$$

The solution of (14) is expressed in terms of the solution  $y(\alpha; t)$  of (5) and the solution  $\tilde{y}(\alpha; t)$  of the adjoint problem

$$(15) \quad \Lambda(\alpha; t)\tilde{y}(\alpha; t) = 0, \quad \tilde{y}(\alpha; T) = 0, \quad (d/dt)\tilde{y}(\alpha; T) = 1.$$

Applying the method of variation of the parameters, one can prove that the solution of (14) is given by

$$(16) \quad z(\alpha; t) = \left[ y(\alpha; t) \int_0^t y(\alpha; t')\tilde{y}(\alpha; t')q(t') dt' - \tilde{y}(\alpha; t) \int_0^t y(\alpha; t')y(\alpha; t')q(t') dt' \right] / W,$$

where

$$(17) \quad W = p(t)[y d\tilde{y}/dt - \tilde{y} dy/dt] = \text{const.}$$

Hence

$$(18) \quad z(\alpha; T) = \frac{dy(\alpha; T)}{d\alpha} = y(\alpha; T) \int_0^T \tilde{y}(\alpha; t)y(\alpha; t)q(t) dt / W = y(\alpha; T) \int_0^T G_{\lambda=0}(t, t, \alpha)q(t) dt$$

where  $G_{\lambda}(t, t', \alpha)$  is again Green's function for the boundary-value problem (4), expressed in terms of the solutions of the initial-value problem (5) and its adjoint (15). A comparison of (18) with (12) completes the proof.

We now turn to the generalization of Theorem 1 to the  $N$  dimensional case. The functions  $p(t)$  and  $q(t)$  of (3) are now replaced by the  $N \times N$  symmetric matrices  $P(t)$  and  $Q(t)$ .  $P(t)$  is required to be positive definite for all  $0 \leq t \leq T$ . The  $N$  dimensional vectors on which  $\Lambda(\alpha; t)$  acts are denoted by  $\dot{U}(\alpha; t) \equiv (U_1(\alpha; t), \dots, U_N(\alpha; t))$ .

**THEOREM 2.** *Let a differential operator*

$$(19) \quad \Lambda(\alpha; t) = (d/dt)(P(t)d/dt) + \alpha Q(t)$$

be defined for  $0 \leq t \leq T$ ,  $0 \leq \alpha \leq 1$ , with  $P(t)$  and  $Q(t)$  as discussed above.

Let  $\lambda_k(\alpha)$  and  $\vec{U}^{(k)}(\alpha; t)$  be the eigenvalues and eigenfunctions for the boundary-value problem

$$(20) \quad \Lambda(\alpha; t)\vec{U}(\alpha; t) + \lambda(\alpha)\vec{U}(\alpha; t) = 0, \quad \vec{U}(\alpha; 0) = \vec{U}(\alpha; T) = 0.$$

Let  $\vec{y}^{(j)}(\alpha; t)$  ( $j = 1, \dots, N$ ) be the  $N$  independent solutions of the initial-value problem

$$(21) \quad \Lambda(\alpha; t)\vec{y}^{(j)}(\alpha; t) = 0, \quad \vec{y}^{(j)}(\alpha; 0) = 0, \quad (d/dt)y_i^{(j)}(\alpha; 0) = \delta_{ij}.$$

Let  $D(\alpha)$  be

$$(22) \quad D(\alpha) = \det[y_i^{(j)}(\alpha; T)].$$

Then, for every  $\alpha$

$$(23) \quad \left| \prod_{k=1}^{\infty} \frac{\lambda_k(\alpha)}{\lambda_k(0)} \right| = \left| \frac{D(\alpha)}{D(0)} \right|.$$

OUTLINE OF THE PROOF. The proof of Theorem 2 is again based on evaluating the logarithmic derivative of the two sides of (23). Once again it is shown that the logarithmic derivative can be expanded in the form

$$(24) \quad \frac{df(\alpha)}{d\alpha} / f(\alpha) = \sum_{j,i=1}^N \int_0^T dt G_{j,i}^{(\lambda=0)}(t, t, \alpha) Q_{ij}(t) = \frac{d\tilde{f}(\alpha)}{d\alpha} / \tilde{f}(\alpha)$$

where  $f(\alpha)$  and  $\tilde{f}(\alpha)$  denote the functions of  $\alpha$  on the left- and right-hand sides of equation (23) respectively. The Green's "function" is a matrix  $G_{ij}^{(\lambda)}(t, t', \alpha)$ . In proving (24) use is made of the two equivalent methods to express the Green's function, the spectral representation and the representation by means of the  $2N$  independent solutions of equation (21) and its adjoint.

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#### REFERENCES

1. R. P. Feynman and A. R. Hibbs, *Quantum mechanics and path integrals*, McGraw-Hill, New York, 1965.
2. I. M. Gel'fand and A. M. Yaglom, *J. Math. Phys.* **1** (1960), p. 48.
3. M. C. Gutzwiller, *J. Math. Phys.* **8** (1967), p. 1979.
4. S. Levit and U. Smilansky, *Ann. Phys.* **103** (1977), p. 198.
5. P. M. Morse and H. Feshbach, *Methods of theoretical physics. I*, McGraw-Hill, New York, 1953, pp. 1-997. MR **15**, 583.
6. H. Poincaré, *Acta Math.* **4** (1884), p. 213.
7. M. Morse, *The calculus of variations in the large*, Amer. Math. Soc. Colloq. Publ., vol. 18, Amer. Math. Soc., Providence, R. I., 1934; reprint 1966.

DEPARTMENT OF NUCLEAR PHYSICS, WEIZMANN INSTITUTE OF SCIENCE, REHOVOT, ISRAEL