

Concepts of condensed matter physics - Exercise #6

Spring 2018

Due date: 05/07/2018

- 1. Quantization of inter-layer Hall conductivity** – Consider two parallel two-dimensional conductors (e.g. two layers of graphene separated by a thin insulating buffer). In the presence of a perpendicular magnetic field the generalized conductivity tensor has the following form

$$J_i^a = \sigma_{ij}^{ab} E_j^b$$

where $a = 1, 2$ denotes the layer number and J_i^a and E_i^a are the i 'th component of the in-plane current density and electric field. Show that (also) the off-diagonal Hall conductivity σ_{xy}^{12} is quantized following the arguments presented in class. Do so following these steps:

- Write the generalized linear response formula for this off diagonal element, $\sigma_{xy}^{12} = \frac{J_x^1}{\epsilon_y^2}$.
 - Using the appropriate toroidal geometry, write this element in terms of derivatives of the ground state with respect to magnetic fluxes threaded through the torus' holes.
 - Generalize the argument given in class to prove that σ_{xy}^{12} is quantized. What happens when the two layers are decoupled?
- 2. A simple tight-binding model for a 2D topological insulator** - Discuss σ_{xy} for spin-less particles on a square lattice model that has the following Hamiltonian

$$H = \sum_{\mathbf{k}} (\psi_s^+(\mathbf{k}) \quad \psi_p^+(\mathbf{k})) \hat{H}(\mathbf{k}) \begin{pmatrix} \psi_s(\mathbf{k}) \\ \psi_p(\mathbf{k}) \end{pmatrix}, \text{ where}$$

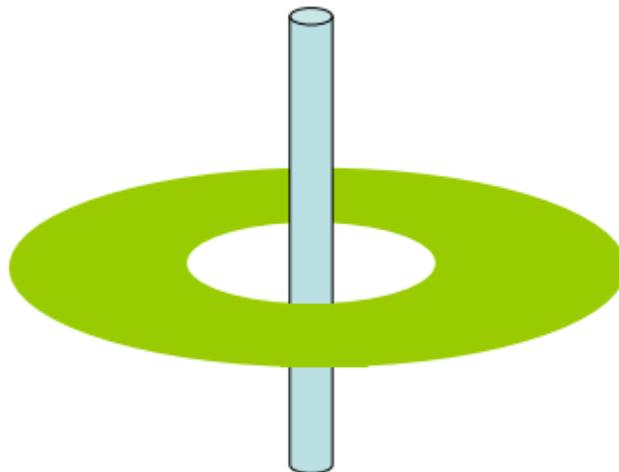
$$\hat{H}(\mathbf{k}) = A(\sin k_x \tau_x + \sin k_y \tau_y) + (m - t \cos k_x - t \cos k_y) \tau_z.$$

Here the τ 's are Pauli matrices acting in the orbital basis.

- Find the corresponding real-space representation of the tight-binding Hamiltonian.
- Discuss σ_{xy} as a function of m

- c. Plot the pseudo spin configuration as a function of $e = m/t$ for different values of e_a -- choose them wisely.
- d. Assume that the crystal exists only for $x < 0$, and that for $x > 0$ there is vacuum. Write the Schrodinger equation for the single particle solutions near the Fermi energy and (assume that $m > 0$ and that e is close to the critical value).
 - i. What are the boundary conditions at $x = 0$?
 - ii. What are the conditions for the existence of a gapless solution on the boundary?
 - iii. What is the decay length of the wave function?
 - iv. What happens to the solution at the critical value of the parameter e ?
- e. Now assume that the crystal exists for all x . Consider the situation where for $x < 0$ the parameter e is slightly larger than the critical value, and for $x > 0$ the parameter e is slightly smaller than it. Find the gapless 1D mode residing on the boundary.
- f. Can you generalize the model to one that realizes an arbitrary Chern number?

3. Laughlin's argument and a preview to the fractional quantum Hall effect – Consider a quantum Hall state on an annulus, as shown in the figure below.



Imagine threading magnetic flux through the hole.

- a. Show, using classical electrodynamics, that a charge flows from the inner edge to the outer edge as a result of changing the flux.
- b. Consider the situation where the flux is increased very slowly from 0 to ϕ_0 . Relate the total charge transferred between the edges during the process to the Hall conductance σ_{xy} .
- c. Use the above argument and the known properties of the Landau levels to deduce σ_{xy} in cases where an integer number of Landau levels are filled (neglecting interactions). What can you say about the robustness of these results in the presence of interactions?
- d. What is the charge that moved from the interior to the exterior if $\sigma_{xy} = \frac{e^2}{3h}$ (this situation corresponds to the $\nu = \frac{1}{3}$ fractional quantum Hall state, observed in experiments). Use the previous sections, and the adiabatic theorem to deduce that the quasiparticles carry fractional charges.