## Home Work 3

- 1. (Ch. 5, ex. 5.1) Consider a particle in a harmonic well, starting at t=0 with  $q=-q_0$  and p=0. Sketch the location of the particle in phase space at  $t=0, 1/4\tau, 1/2\tau, 3/4\tau$ , and  $\tau$ , where  $\tau$  is a period of motion. Connect the points and show that the orbit is an ellipse. Which way does the phase space orbit circulate, clockwise or counterclockwise?
- 2. (Ch. 5, ex. 5.2) Consider a pendulum, starting at t=0 with q=0 and  $p=p_0$ . Sketch the location of the particle in phase space at t=0,  $1/4\tau$ ,  $1/2\tau$ ,  $3/4\tau$ , and  $\tau$ , where  $\tau$  is a period of motion. You should get discover three different regions, two corresponding to rotation and one to libration, depending on the value of  $p_0$ . Show that the direction of the rotational orbits in phase space correlate smoothly with those of the libration at separatrix.
- 3. (Ch. 5, ex. 5.3) To understand why the Fourier transform variable p in the definition of the Wigner function has the physical meaning of a momentum, examine the structure of the quantity  $\langle x'|\Psi\rangle\langle\Psi|x\rangle$  for the case where  $\Psi(x)=Ne^{-\alpha(x-x_0)^2+ip_0(x-x_0)}$ . This product is a function of the variables x and x' and is complex. Plot the real part of this product and discuss the wavelength of oscillations along the s=x-x' coordinate. If the object is Fourier transformed along the s=x-x' coordinate, where will its peak be in the Fourier variable, p?
- 4. (Ch. 5, ex. 5.4) The Wigner distribution has several important properties which suggest its interpretation as a probability distribution:

$$\int_{-\infty}^{\infty} f_W(p,q) \, dp = \langle q | \Psi \rangle \langle \Psi | q \rangle = |\Psi(q)|^2$$
 (4.1)

$$\int_{-\infty}^{\infty} f_W(p,q) \, dq = \langle p | \Psi \rangle \langle \Psi | p \rangle = |\tilde{\Psi}(p)|^2$$
 (4.2)

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f_W(p, q) \, dp dq = 1 \tag{4.3}$$

where  $\Psi(p)$  is the momentum representation of the wavefunction  $|\Psi\rangle$ . Derive these relations.

5. (Ch. 5, ex. 5.5) (divided here into a preliminary and a full calculation)) Consider the normalized wavepacket

$$\Psi(q,t) = \left(\frac{2Re\alpha_t}{\pi}\right)^{1/4} e^{-\alpha_t (q-q_t)^2 + \frac{i}{\hbar} p_t (q-q_t) + \frac{i}{\hbar} \gamma_t},\tag{5.1}$$

and assume that  $\gamma_t$  is real.

a. Taking  $\alpha_t = \frac{m\omega}{2\hbar}$ , show that

$$f_W(q,p) = \frac{1}{\pi \hbar} e^{-\frac{m\omega}{\hbar} (q - q_t)^2} e^{-\frac{1}{m\omega\hbar} (p - p_t)^2}.$$
 (5.2)

b. Show that in the more general case, where

$$\alpha_t = a(\frac{\alpha_0 \cos \omega t + ia \sin \omega t}{i\alpha_0 \sin \omega t + a \cos \omega t}) \tag{5.3}$$

 $(a = \frac{m\omega}{2\hbar})$ , the Wigner distribution takes the form:

$$f_W(q,p) = \frac{1}{\pi \hbar} e^{-\frac{2|\alpha_t|^2}{\text{Re }\alpha_t}(q-q_t)^2} e^{-\frac{1}{2\hbar^2 \text{Re }\alpha_t}(p-p_t)^2} e^{-\frac{2\text{Im }\alpha_t}{\hbar \text{Re }\alpha_t}(q-q_t)(p-p_t)}.$$
 (5.4)

6. (Ch. 5, EOC ex. 1) Calculate  $Tr(\rho)$  and  $Tr(\rho^2)$  for the generic 2x2 density matrix:

$$\rho_{mn} = \begin{pmatrix} |a|^2 & a^*b \\ ab^* & |b|^2 \end{pmatrix} \tag{6.1}$$

Then repeat, this time with the off-diagonal elements set equal to 0. Show that in the first case  $Tr(\rho^2) = 1$  (pure state) and in the second case  $Tr(\rho^2) \le 1$  (mixed state).

7. (Ch. 5, EOC ex. 3) Verify that

$$Tr(
ho^2) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dp \int_{-\infty}^{\infty} dq 
ho_W^2(p,q)$$

for the harmonic oscillator potential. First calculate the LHS using both the energy representation for  $\rho$ :

$$\langle \psi_n | \rho | \psi_m \rangle = e^{-\beta E_n} / Q \delta_{nm} \qquad Q = \sum_n e^{-\beta E_n}.$$
 (7.1)

Then repeat the calculation using the x representation for  $\rho$ :

$$\frac{\langle x'|e^{-\beta H}|x\rangle}{Q} = \left(\frac{m\omega}{\pi\hbar}\tanh(f/2)\right)^{1/2} exp\left\{\frac{-m\omega}{2\hbar\sinh(f)}[(x^2 + x'^2)\cosh(f) - 2xx']\right\}$$
(7.2)

where  $f = \hbar \omega/kT$  and  $Q = \int_{-\infty}^{\infty} \langle x|e^{-\beta H}|x \rangle dx$ . Note that the higher the temperature, the farther  $\rho$  gets from a pure state, i.e.  $Tr(\rho^2)$  decreases with temperature. Plot  $Tr(\rho^2)$  as a function of T.