

# Resonant Eigenfunction **Delocalization** of Random Schrödinger Operators on Tree Graph

1. *Unbounded potentials: **Extended States in a Lifshitz Tail Regime***
2. *Bounded potentials: **Absence of Mobility Edge** at Weak Disorder*

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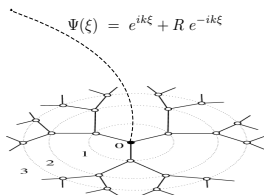
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# The random Hamiltonian

$$H_\lambda(\omega) := T + \lambda V(x; \omega)$$

- on the homogeneous tree graph of degree  $K + 1$ ; **Bethe lattice**  $\mathbb{B}$ .



Hopping term:  $(T\psi)(x) := -\sum_{\text{dist}(x,y)=1} \psi(y)$

Random potential:  $V(x; \cdot)$ ,  $x \in \mathbb{B}$ , i.i.d. random variables,

$\mathbb{P}(V(0) \in dv) = \varrho(v) dv$ ,  $\varrho$  bounded, piecewise monotone & piecewise cont.

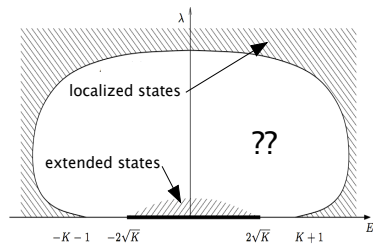
In this talk, special emphasis on two cases:

- $\text{supp } \varrho = \mathbb{R}$  (e.g.  $V(x)$  Gaussian or Cauchy)
- $\text{supp } \varrho = [-1, 1]$  (bounded potential,  $|V(x; \omega)| \leq 1$ )

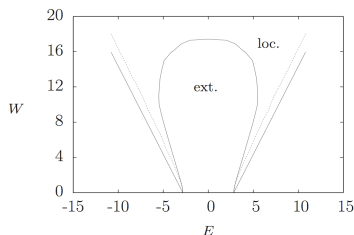
Alternative form:

$$\mathcal{H}(\omega) = -t \sum_{\langle x,y \rangle \subset \mathbb{B}} a_x^\dagger a_y + \frac{W}{2} \sum_{x \in \mathbb{B}} V(x; \omega) a_x^\dagger a_x$$

# Previous results and expected phase diagrams



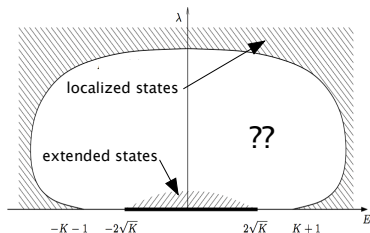
phase diagram for  
unbounded potential



expected mobility edge  
for a bounded potential

- Among the earliest studied models of the **Anderson localization** And. '58  
Abou-Chacra/Anderson/Thouless '73, Abou-Chacra/Thouless '74
- **Motivation:** Relatively more accessible compared to  $\mathbb{Z}^d$ .  
Self-consistent approach to localization becomes exact ( $\neq$  solvable !).
- Renewed interest as a model for the configuration space of systems with **many particles** Altshuler/Gefen/Kamenev/Levitov '97, (cf. Pal/Huse'11)
- **Numerical work:** Miller/Derrida '94, Biroli/Semerjian/Tarzia '10
- **Rigorous results:** *next page*

# Some previous rigorous results



$$H = T + \lambda V$$

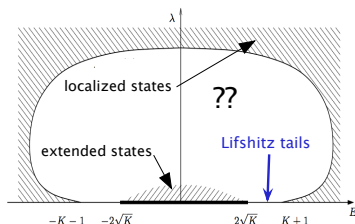
supp  $\varrho = \mathbb{R}$

+ certain regularity conditions

- Spectrum of the Laplacian on  $\ell^2(\mathbb{B})$ :  $\sigma(T) = [-2\sqrt{K}, 2\sqrt{K}]$

- 1 **pp (loc.) spectrum** at strong disorder: Aizenman/Molchanov '93  
Aiz. '94  
and at large energies
- 2 **ac (ext.) spectrum** for weak disorder at energies within  $\sigma(T)$  Klein '94  
Aiz./Sims/Warzel '05, Froese/Hasler/Spitzer '06
- 3 **Lifshitz tails**: For  $\lambda \rightarrow 0$ , the mean density of states ( $\rho_{\text{DOS}}$ ) at energies outside  $\sigma(T)$  **vanishes faster than any power of  $\lambda$**  Miller/Derrida '94

# The long open puzzle



$$\text{supp } \varrho = \mathbb{R}$$

$$\int_{\mathbb{R}} v \varrho(v) dv = 0$$

**Question:** Where is the edge of the localization regime, in particular, at **weak disorder**?

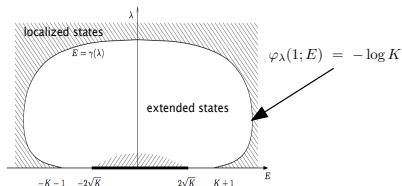
**Note:** for energies  $E$  outside  $\sigma(T) = [-2\sqrt{K}, 2\sqrt{K}]$ , the **mDOS** ( $\rho_{\text{DOS}}$ ) vanishes at weak disorder to all orders in perturbation theory (**Lifshitz tail regime**)

E.g., in case of the Gaussian random potential, for  $E \notin \sigma(T)$ :

$$\rho_{\text{DOS}}(E) \approx \exp\left(-C(E)/\lambda^2\right) \quad \text{as } \lambda \downarrow 0.$$

## A somewhat surprising answer

**Theorem:** In the case of unbounded random potential ( $\text{supp } \rho = \mathbb{R}$ , etc.), for  $\lambda > 0$  the **ac spectrum** immediately extends up to  $E = \pm(K + 1)$ , in particular, into the regime of **Lifshitz tails**.



More can be said in terms of the Green function

$$G(0, x; E) := \left\langle \delta_0, (H - E - i0)^{-1} \delta_x \right\rangle$$

and its **moment generating function**

$$\text{for } s < 1 : \quad \varphi_\lambda(s; E) := \lim_{|x| \rightarrow \infty} \frac{\log \mathbb{E} [ |G_\lambda(0, x; E + i0)|^s ]}{|x|}$$

$$\text{for } s = 1 : \quad \varphi_\lambda(1; E) := \lim_{s \nearrow 1} \varphi_\lambda(s; E)$$

The mechanism at work here: fluctuation enabled resonant tunneling:

States which locally appear to be localized have **arbitrarily close energy gaps** ( $\Delta E$ ) with other states (**at distances**  $R$ ), to which the **tunneling amplitudes** are **exponentially small** (as  $\approx e^{-L_\lambda(E)R}$ ).

**Mixing** between two levels occurs if

$$\Delta E \ll e^{-L_\lambda(E)R}$$

Since the **volume grows exponentially fast** (as  $K^R$ ),  
extended states will form provided

$$L_\lambda(E) < \log K$$

Essential enabling conditions:

- local fluctuations in the self energy
- the exponential growth of the configuration space volume

## (Almost) complementary criteria for pp and ac spectra

Assumptions:  $\varrho(V) > 0$  on  $\mathbb{R}$ ,  $\int |v|^\tau \varrho(v) dv < \infty$  for some  $\tau > 0$ , etc.

Theorem 1 (localization [Aiz./Molchanov '93, Aiz. '94])

- If for all (or a.e.) energies  $E$  in some interval  $I \subset \mathbb{R}$

$$\varphi_\lambda(1; E) < -\log K \quad (1)$$

then  $H(\omega)$  has only pure point (localized) spectrum in that interval.

- Furthermore, at weak disorder (1) holds for energies  $|E| > (K + 1)$ .

The new, *complementary*, statement:

Theorem 2 (delocalization [Aiz./Warzel '10, '11])

- Under the above assumptions on  $\rho$ , at energies at which

$$\varphi_\lambda(1; E) > -\log K \quad (2)$$

one has:  $\operatorname{Im} G(x, x; E + i0) > 0;$

- ( $\implies$ ) if (2) holds for a positive measure of energies  $E \in I$ , then  $H(\omega)$  has absolutely continuous (delocalized) spectrum in that interval.

Added by M. Shamis: One may conclude from Thm. 2 that if (2) holds for almost every  $E \in I$  then then  $H(\omega)$  has only ac spectrum in  $I$ .



# Implications for the phase diagram

Assuming:

- i. the function  $\varphi_\lambda(1; E)$ , and/or just the related **Lyapunov exponent**, are **continuous in  $(E, \lambda)$**  (there are gaps here in the mathematical theory),
- ii. the equality  $\varphi_\lambda(1; E) = -\log K$  holds only along a simple curve,

one may conclude:

1. the random operator has a **(simple) mobility edge** which converges to  $|\mathbb{E}| = \pm(K + 1)$  **(and not to the edge of the  $\lambda = 0$  spectrum  $\sigma(T)$ )**.
2. up to the mobility edge the random operator has **(a.s.) purely extended states**; beyond it only pure point spectrum, with dynamical localization and all that.

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For **bounded potentials**, Theorem 2 yields another surprise:

**There is no such mobility edge at weak disorder !**

the (generally) expected phase diagram needs to be corrected.



References (to the new results presented here):

1. M. Aizenman, S. Warzel, "Extended States in a Lifshitz Tail Regime for Random Schrödinger Operators on Trees", *Phys. Rev. Lett.* **106**, 136804 (2011).
2. M. Aizenman, S. Warzel, "Resonant delocalization for random Schrödinger operators on tree graphs" <http://arxiv.org/abs/1104.0969>.
3. M. Aizenman, S. Warzel, " Absence of Mobility Edge for the Anderson Random Potential on Tree Graphs at Weak Disorder, ([in preparation, to be submitted](#)).

