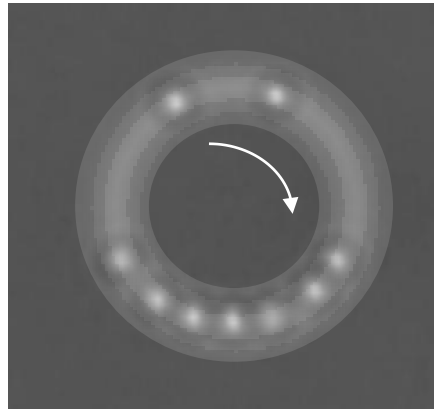


Non-equilibrium pair attraction between driven colloidal particles

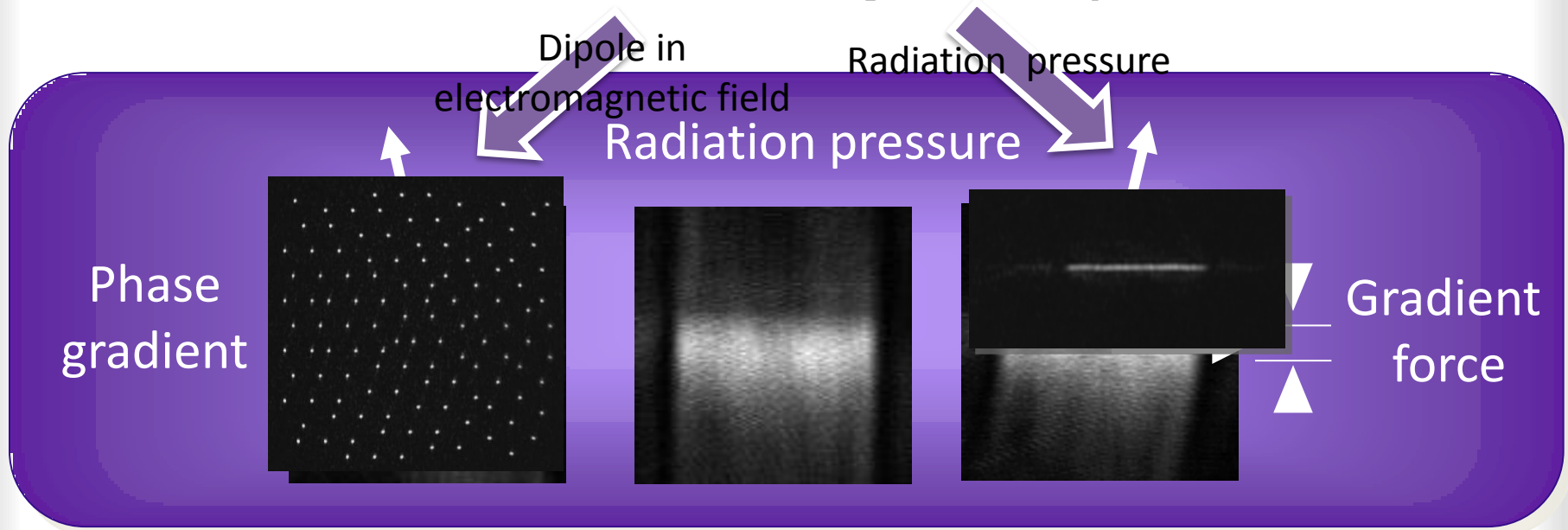


Yulia sokolov, Derek Frydel,
Haim Diamant and Yael Roichman

Optical trapping

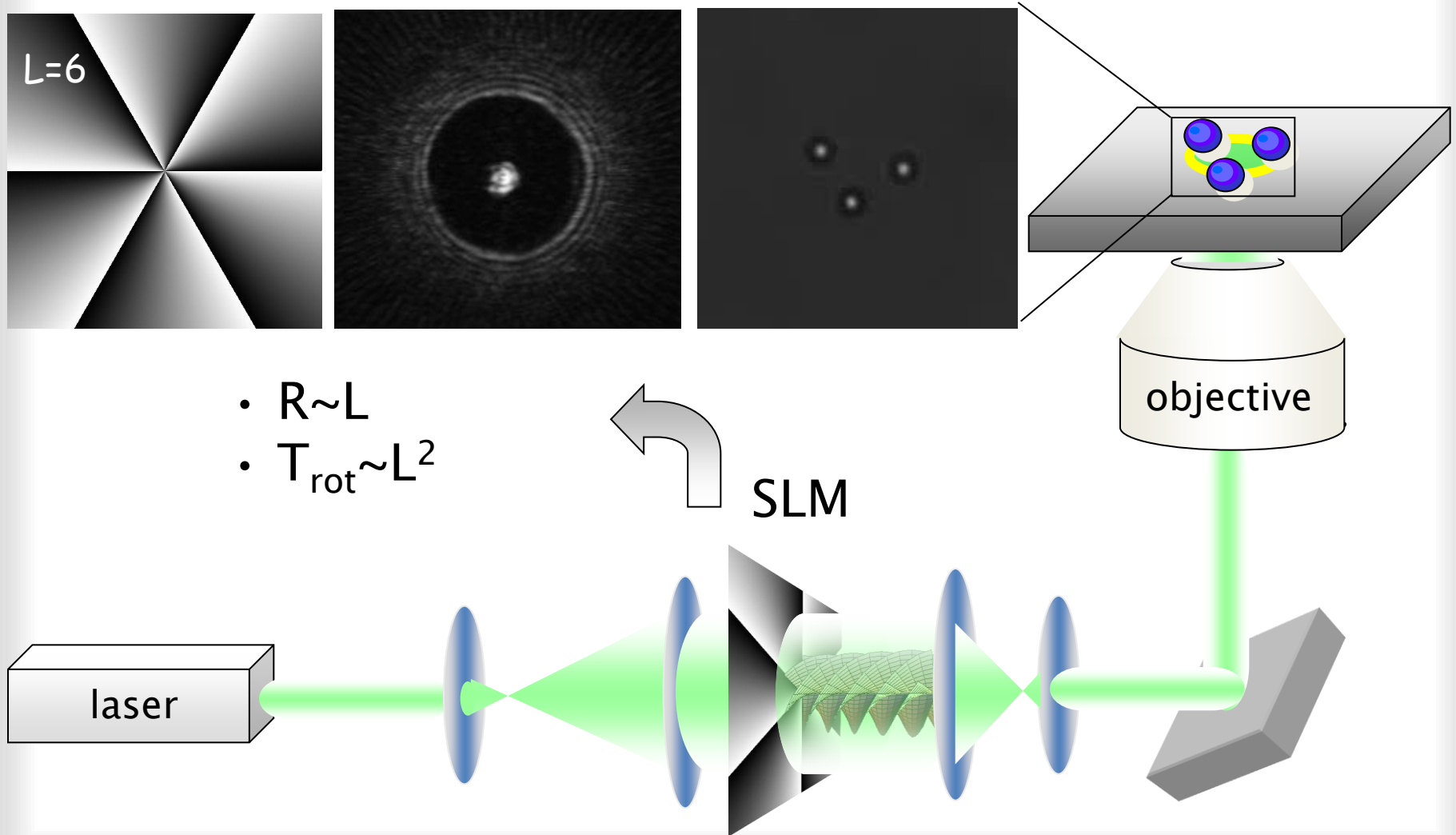


$$\vec{F} = \alpha \nabla \vec{I} + \beta I \nabla \phi$$



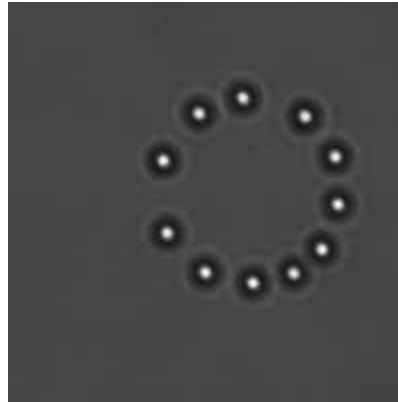
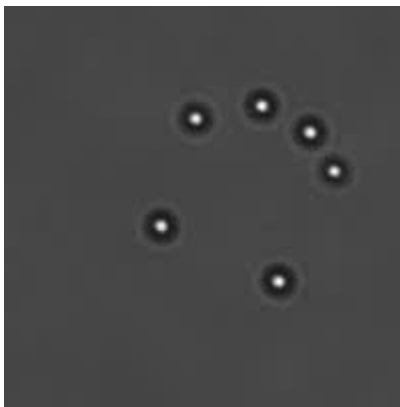
Holographic optical tweezers

Static traps



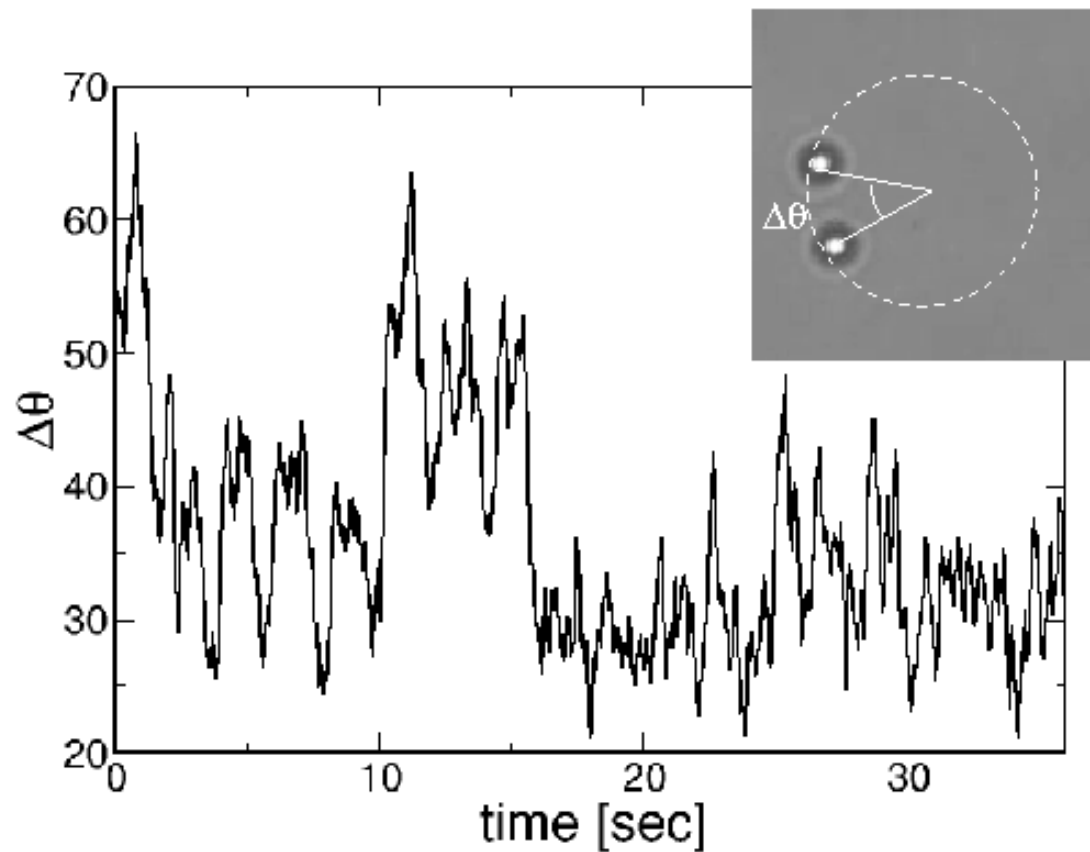
Many particles in an optical vortex

Driven dissipative model system - experiments

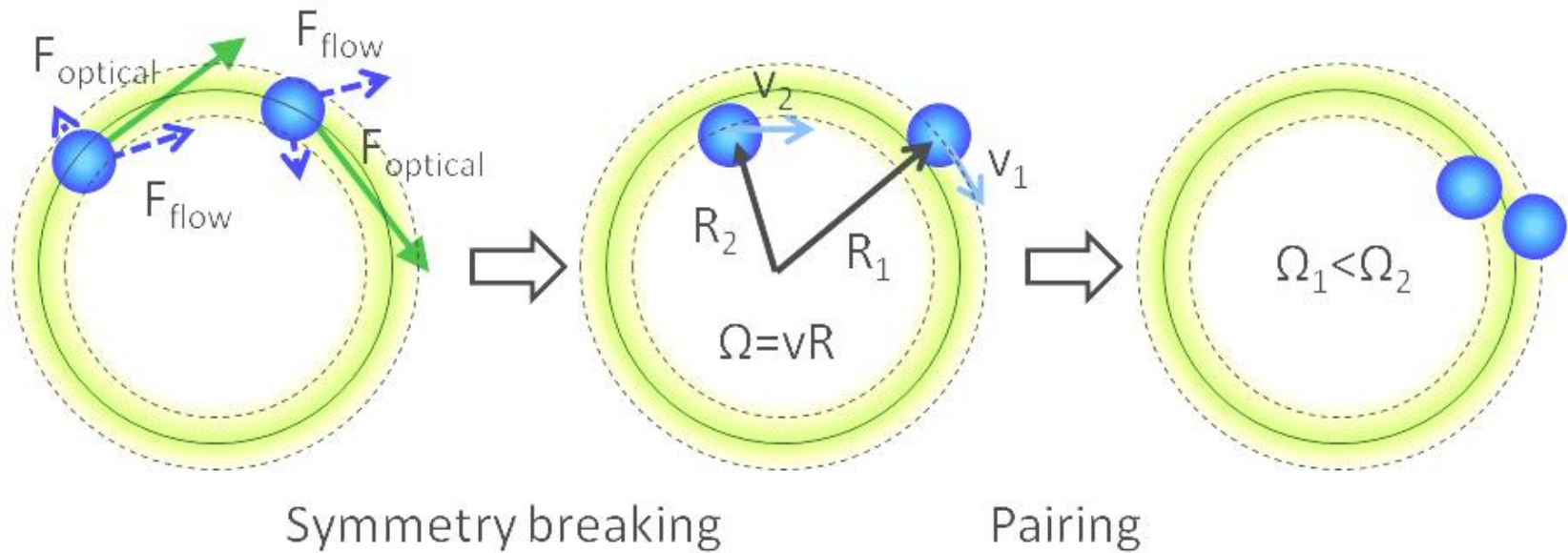


Non-equilibrium, hydrodynamic interaction
phonons, single file diffusion

Two particles - experiments



Two particles - calculation



$$\frac{\gamma \Omega_i}{K_\theta} = 1 + \frac{3a}{16\sqrt{2}} \frac{1}{R} \left[\frac{1 + 3\cos \theta_{ij}}{\sqrt{1 - \cos \theta_{ij}}} \right]$$

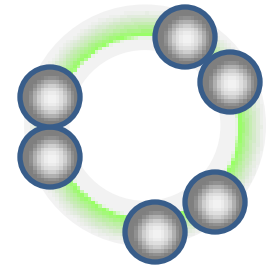
$$\frac{\gamma V_{\rho,i}}{K_\theta} = \frac{3a}{16\sqrt{2}} \left[\frac{3\sin \theta_{ij}}{\sqrt{1 - \cos \theta_{ij}}} \right]$$

Collective mobility calculation

The model:

- Paired particles with “pair” mobility.
- Equidistanced.
- Collective mode.
- Stokeslet approximation.

$$G_{ij} = \frac{1}{8\pi\eta r} \left(\delta_{ij} + \frac{\overline{r_1 r_2}}{r^2} \right)$$



$$\text{collective mobility} = \begin{cases} \text{even } N/2 & \frac{1}{Q\pi\eta r} \left[1 + \frac{Qr}{8\sqrt{2}R} \left(-\frac{1}{\sqrt{2}} + \sum_{n=1}^{(N-4)/4} K(nx) \right) \right] \\ \text{odd } N/2 & \frac{1}{Q\pi\eta r} \left[1 + \frac{Qr}{8\sqrt{2}R} \left(\sum_{n=1}^{(N-2)/4} K(nx) \right) \right] \end{cases}$$

$$K(nx) = \frac{1 + 3\cos(nx)}{\sqrt{1 - \cos(nx)}}$$

$$x = \frac{4\pi}{N}$$

Collective mobility model vs. experiment

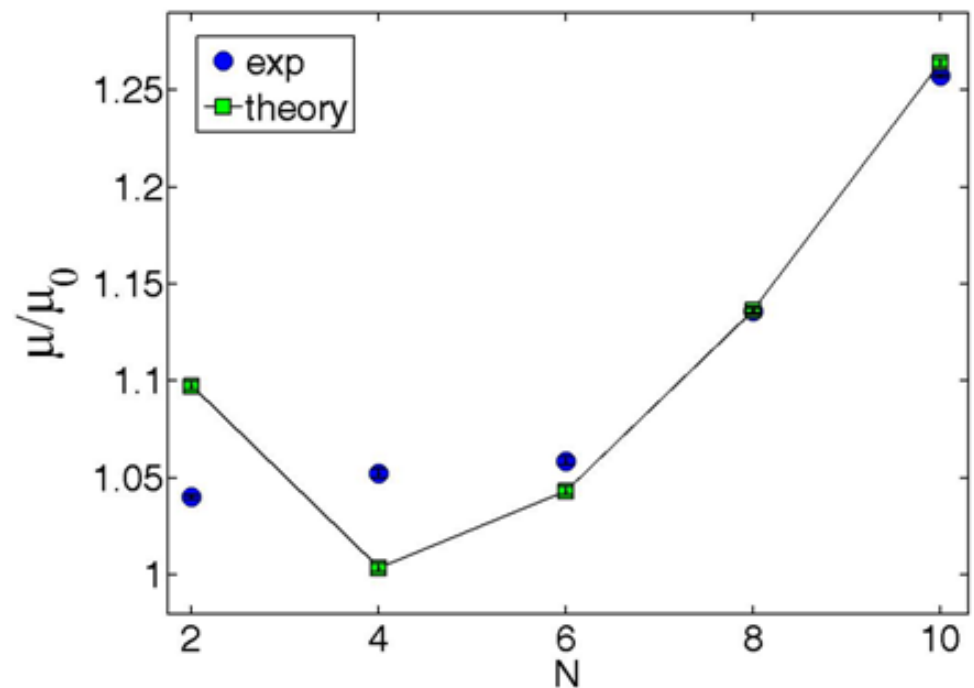
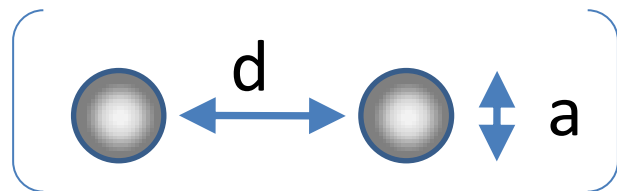
Fitting parameter $\equiv Q$

Drag of a single sphere in fluid

$$\gamma = 6\pi\eta r$$

Drag of coupled spheres in fluid

$$\gamma = Q\pi\eta r$$



Simulation

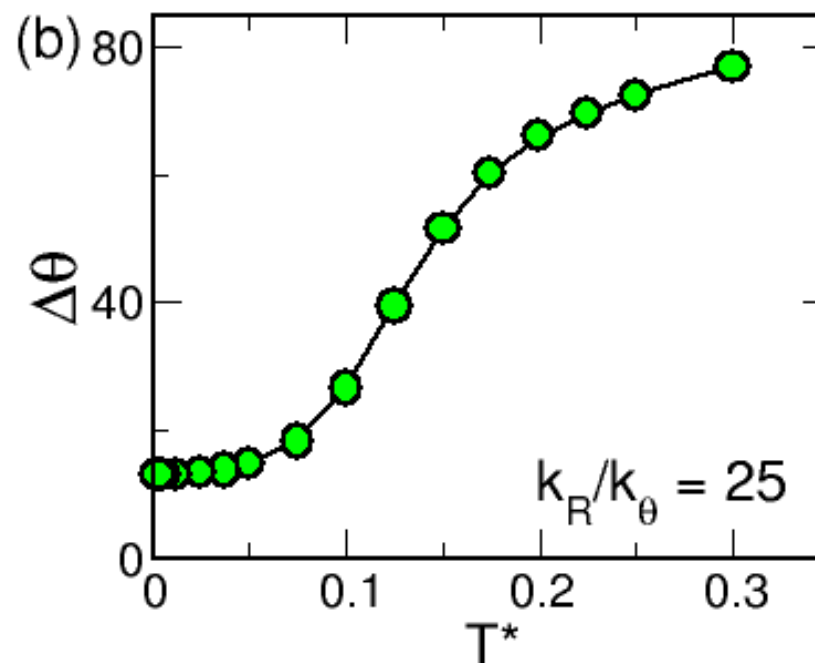
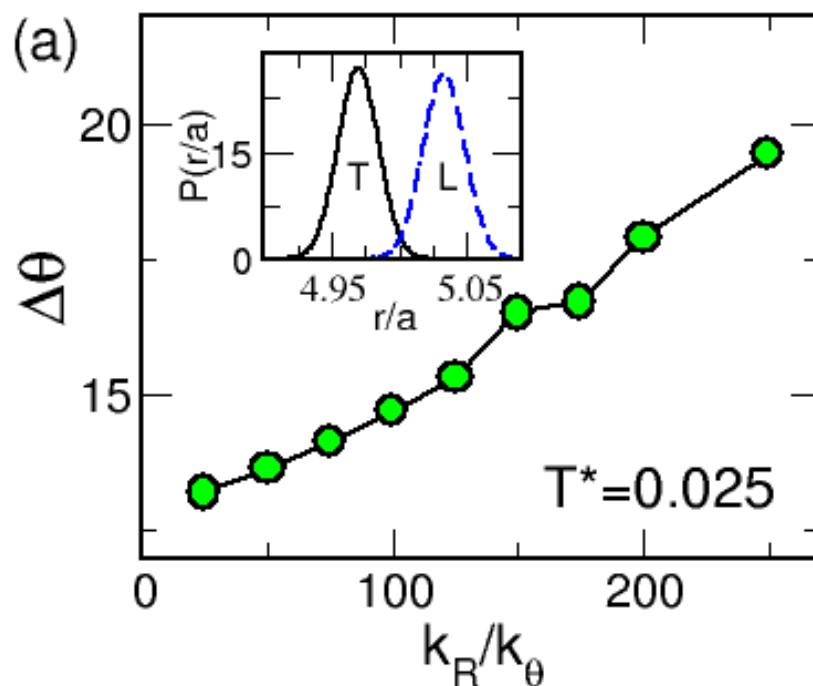
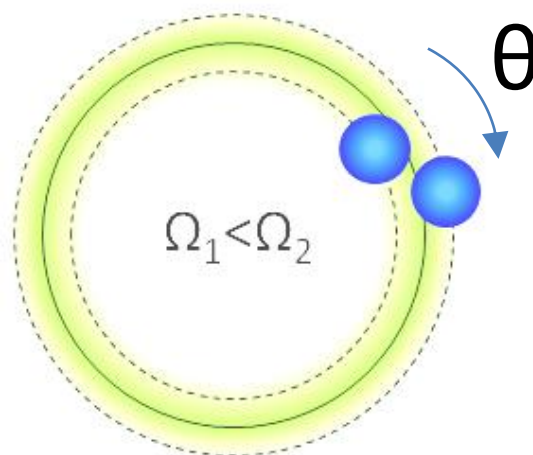
Stokesian dynamics:

- Many body hydrodynamic interactions (Rotne Prager).
- Suitable for systems out of equilibrium.
- Thermal noise included.

We use it to study temperature effects.

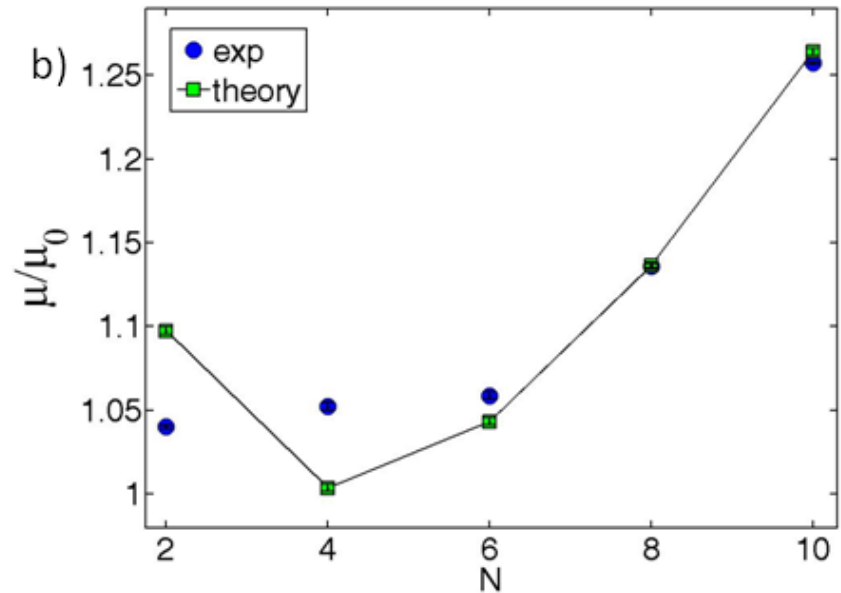
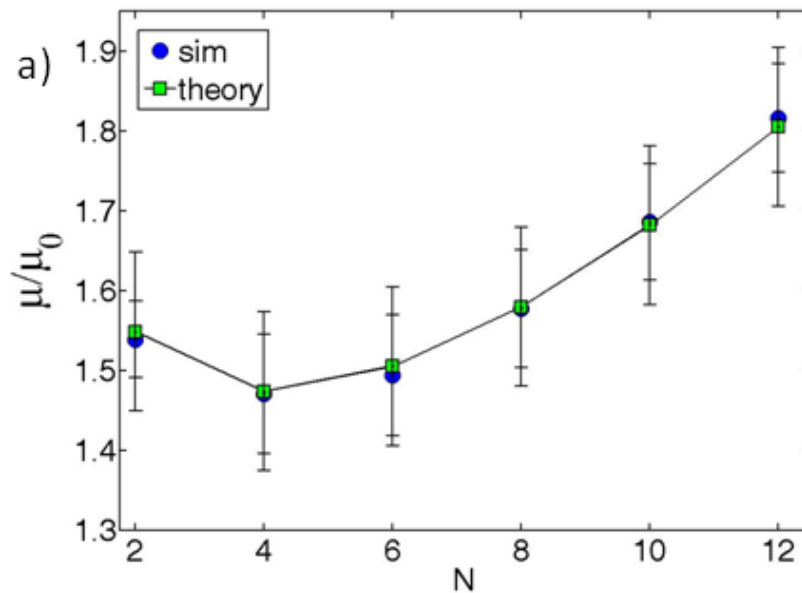
$$m \frac{dv}{dt} = F^{\text{Hydrodynamic}} + F^{\text{Brownian}} + F^{\text{Particles}}$$

Simulations – the origin of pairing



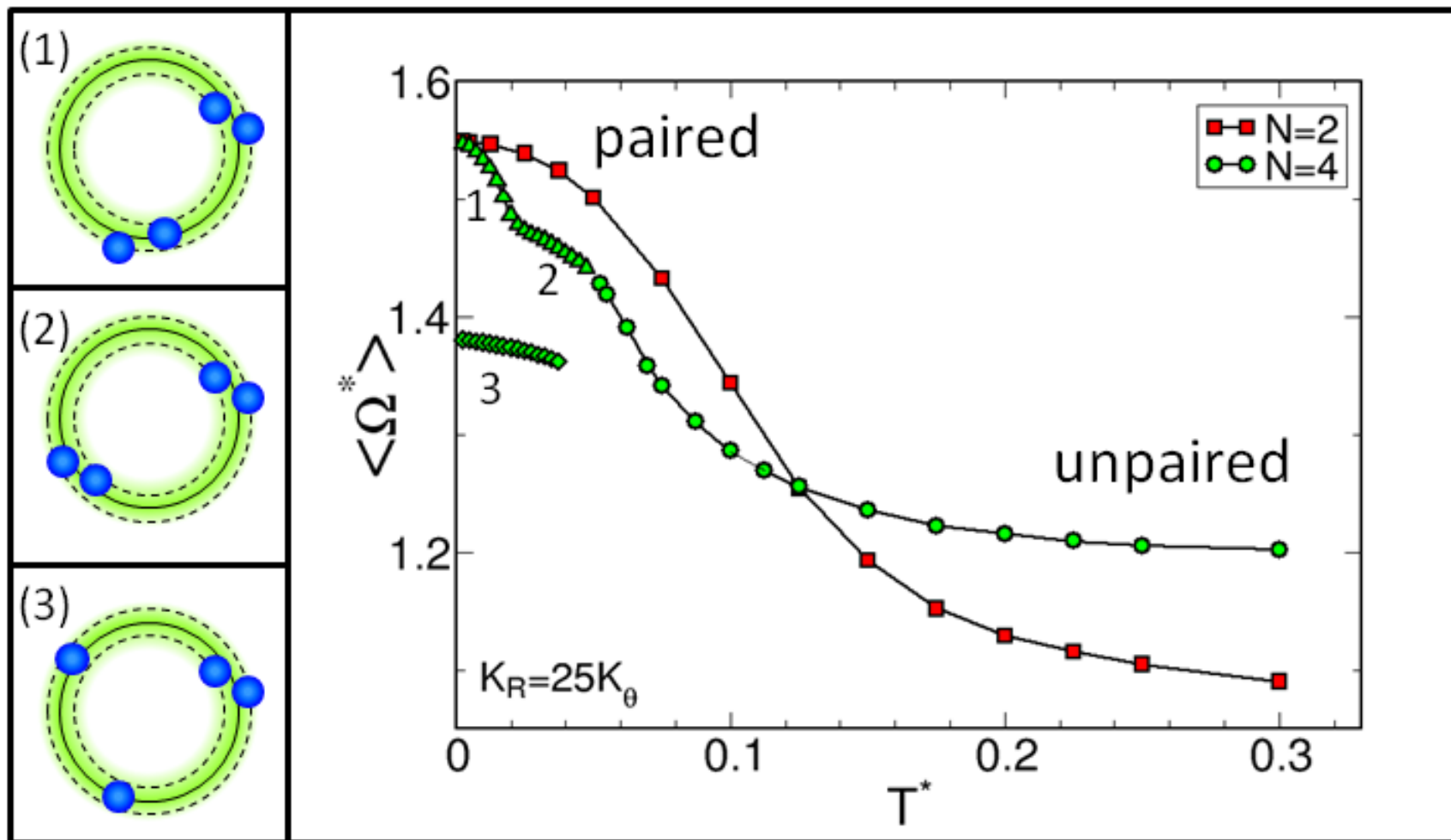
Collective mobility

model, experiment, and simulation



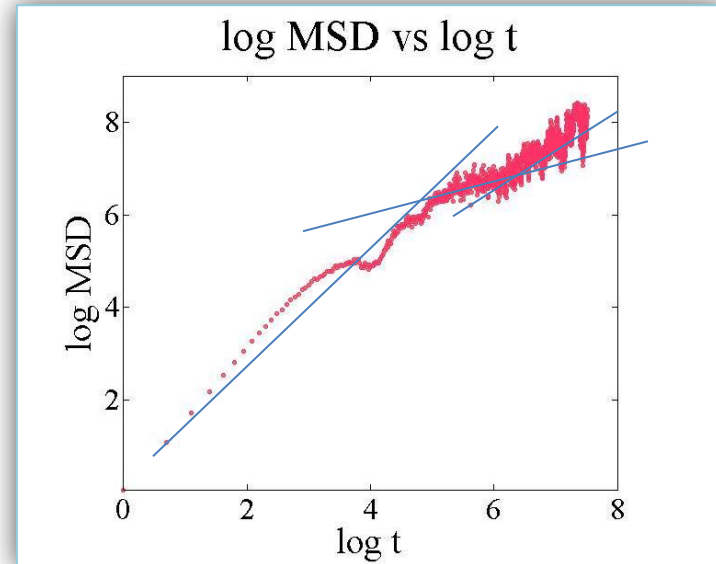
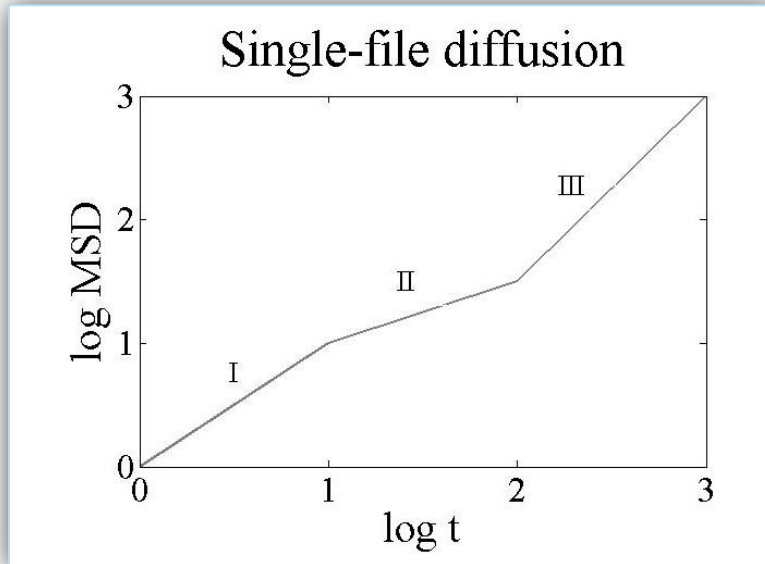
Mean angular velocity

Sensitive macroscopic observable



Single file diffusion

Driven system

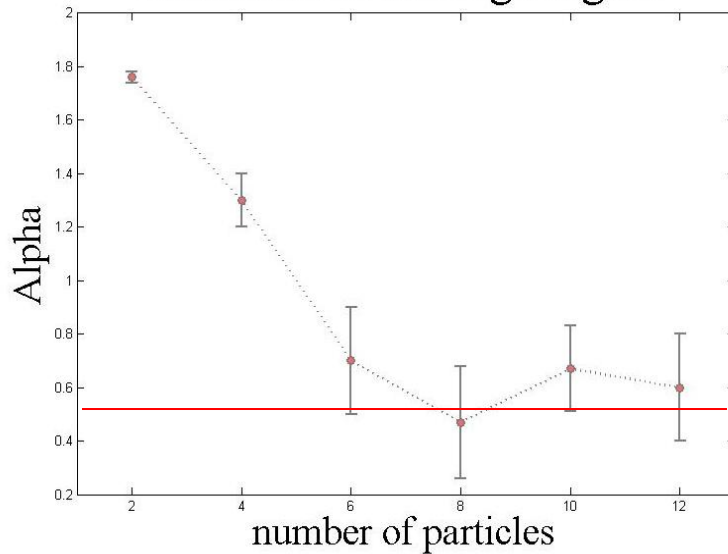


- I – one particle diffusion $\alpha=1$
- II - caging effect, subdiffusion $0.5 < \alpha < 1$
- III - collective diffusion $\alpha=1$

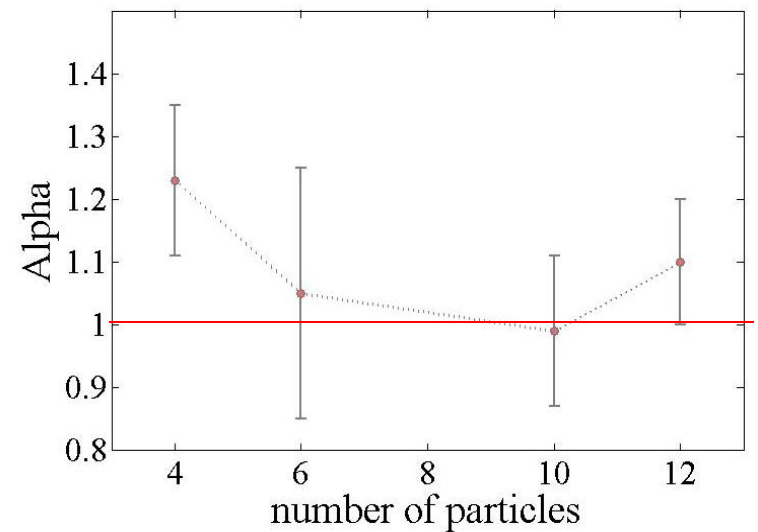
Single file diffusion

Driven system

Diffusion in the cage region



Collective diffusion



Conclusions

Hydrodynamic Interactions give rise to a pseudo-potential in this out-of-equilibrium system, which result in particle pairing.

For an even number of particles there exists a mode in which particles rotate as a “crystal”. Is the system minimizing dissipation while maximizing stability?

Single file diffusion exists in driven systems. We see the crossover to collective diffusion in this small system.