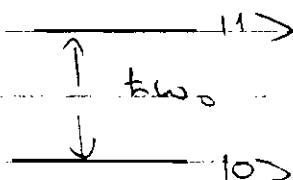


The Qubit:

a two-level quantum system:



a general state:

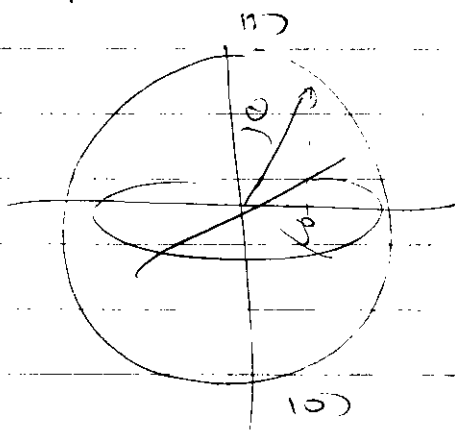
$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle = e^{i\varphi_0} |\alpha| |0\rangle + e^{i\varphi_0} |\beta| |1\rangle$$

$\Rightarrow$ up to a global phase

$$|\alpha|^2 + \text{Re}^2 \beta + \text{Im}^2 \beta = 1$$

2D sphere embedded in 3D:

directions
or Bloch sphere
 $\Rightarrow$ not necessarily
in lab.



Bloch sphere

$\rightarrow$ states that are
orthogonal on Bloch
sphere are not orthogonal
in Hilbert space.

$$\cos\left(\frac{\theta}{2}\right)|1\rangle + e^{i\varphi}\sin\left(\frac{\theta}{2}\right)|0\rangle$$

$$0 < \varphi < 2\pi$$

$$0 < \theta < \pi$$

- when we measure, collapse the qubit onto the measurement basis.

- Amount of info needed to describe 1 qubit = ∞

- Amount of info retrieval after 1 qubit meas = 1 bit.

- to precisely determine the state of 1 qubit

$\Rightarrow \infty$ measurements.

1 copy $\rightarrow \infty$ copies.

- No cloning theorem.

$$U |\psi\rangle |s\rangle \Rightarrow |\psi\rangle |\psi\rangle$$

U is unitary.
clones.

assure for any $|\psi\rangle$ and some $|\psi\rangle$ and $|x\rangle$

$$U |\psi\rangle |\psi\rangle = |\psi\rangle |\psi\rangle$$

$$U |x\rangle |x\rangle = |x\rangle |x\rangle$$

$$\langle \psi | \langle \psi | x \rangle | x \rangle = |\langle \psi | x \rangle|^2$$

can't work for
any state

classically:

$$\langle \psi | \langle \psi | U^\dagger U | x \rangle | \psi \rangle = \langle \psi | x \rangle$$

CNOT

is a clone
 $\Rightarrow$ entangled

$$\Rightarrow \langle \psi | x \rangle = \begin{cases} 1 & \text{either the same} \\ & \text{or orthogonal} \end{cases}$$

operations on a pure state $\Rightarrow$ rotations.

in spin Hilbert space 2×2 SH matrices. $[\sigma_i, \sigma_j] = 2i\epsilon_{ijk} \sigma_k$

3 Generators: $\hat{\sigma}_i = \frac{\hbar}{2} \sigma_i$, $\text{tr}(\sigma_i) = 0$, $\sigma_i \sigma_j = i\epsilon_{ijk} \sigma_k$

$$\sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}; \quad \sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}; \quad \sigma_y = \begin{bmatrix} 0 & i \\ -i & 0 \end{bmatrix}$$

$$R(\theta, \vec{n}) = e^{-i \frac{\theta}{2} \vec{n} \cdot \vec{\sigma}} = \cos\left(\frac{\theta}{2}\right) \hat{1} + i \vec{\sigma} \cdot \vec{n} \sin\left(\frac{\theta}{2}\right)$$

notice that when $\theta = 180^\circ$ $\sin = 1$ $\cos = 0$

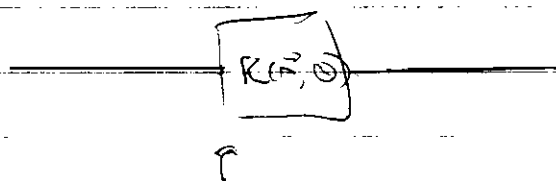
$$\Rightarrow -i \vec{\sigma} \cdot \vec{n} = R(\vec{n}, 180^\circ)$$

σ : Both hermitian and unitary.

- usually in QIP no time evolution $\Rightarrow$ operation.

(no diff eq $\Rightarrow$ linear algebra)

Circuit:



analogy $\Rightarrow$ infinitely many.

Density matrix:

statistical mixture of

states:

$$\hat{\rho} = \sum p_i |\alpha_i\rangle \langle \alpha_i|$$

$$\text{Tr}(\hat{\rho}) = \sum p_i \text{Tr}(|\alpha_i\rangle \langle \alpha_i|) = \sum p_i = 1$$

Similarly:

$$\langle A \rangle = \text{Tr}[\hat{\rho} \hat{A}]$$

By construction $\hat{\rho}$ is hermitian. $p_{ij} = p_{ji}^*$

for a single qubit: # of free parameters:

$$\begin{matrix} \begin{bmatrix} 1 & 2 \\ & 1 \end{bmatrix} & = & 4 - 1 = 3 \\ & \uparrow & \uparrow \\ & \text{hermitian} & \text{Tr}=1 \end{matrix}$$

Since $I, \sigma_x, \sigma_y, \sigma_z$ form a basis for all hermitian

$$\hat{\rho} = \alpha I + \beta \hat{\sigma}_x + \gamma \hat{\sigma}_y + \delta \hat{\sigma}_z$$

$$\begin{aligned} \langle \sigma_x \rangle &= \text{Tr}[\hat{\rho} \sigma_x] = \text{Tr}[\alpha I \sigma_x] + \text{Tr}[\beta \hat{\sigma}_x \sigma_x] + \text{Tr}[\gamma \hat{\sigma}_y \sigma_x] + \text{Tr}[\delta \hat{\sigma}_z \sigma_x] \\ &= 0 + 2\beta + 0 + 0 \\ \beta &= \langle \sigma_x \rangle \frac{1}{2} \end{aligned}$$

$$\Rightarrow \hat{\rho} = \frac{1}{2} \left[\hat{I} + \sum_i \langle \sigma_i \rangle \hat{\sigma}_i \right]$$

state tomography

Geometrically $\hat{\rho}$ is represented by a 3D vector with projections $\langle \sigma_i \rangle$ in the x direction

$$-1 \leq \langle \sigma_i \rangle \leq 1$$

the volume of the Bloch sphere

E.g. fully mixed state $\hat{\rho} = \frac{1}{2} \hat{I}$

$$\langle \sigma_x \rangle = \langle \sigma_y \rangle = \langle \sigma_z \rangle = 0 \quad \text{center of sphere at origin}$$

measure $| \pm i \rangle$ n times.

$$\langle \sigma_i \rangle = \frac{n_+ - n_-}{n}$$

binomial dist.

$$\text{std in } n_+ / n = \sqrt{\frac{pq}{n}}$$

$\langle \sigma \rangle$ known within $\sqrt{\frac{pq}{n}}$

if $p=1, q=0$

$$\sqrt{\frac{pq}{n}} = 0$$

$$p=q=\frac{1}{2} \Rightarrow$$

$$\sqrt{\frac{pq}{n}} = \frac{1}{2\sqrt{n}} \quad 14$$