

Controlled entanglement between levitated masses via parametric resonance

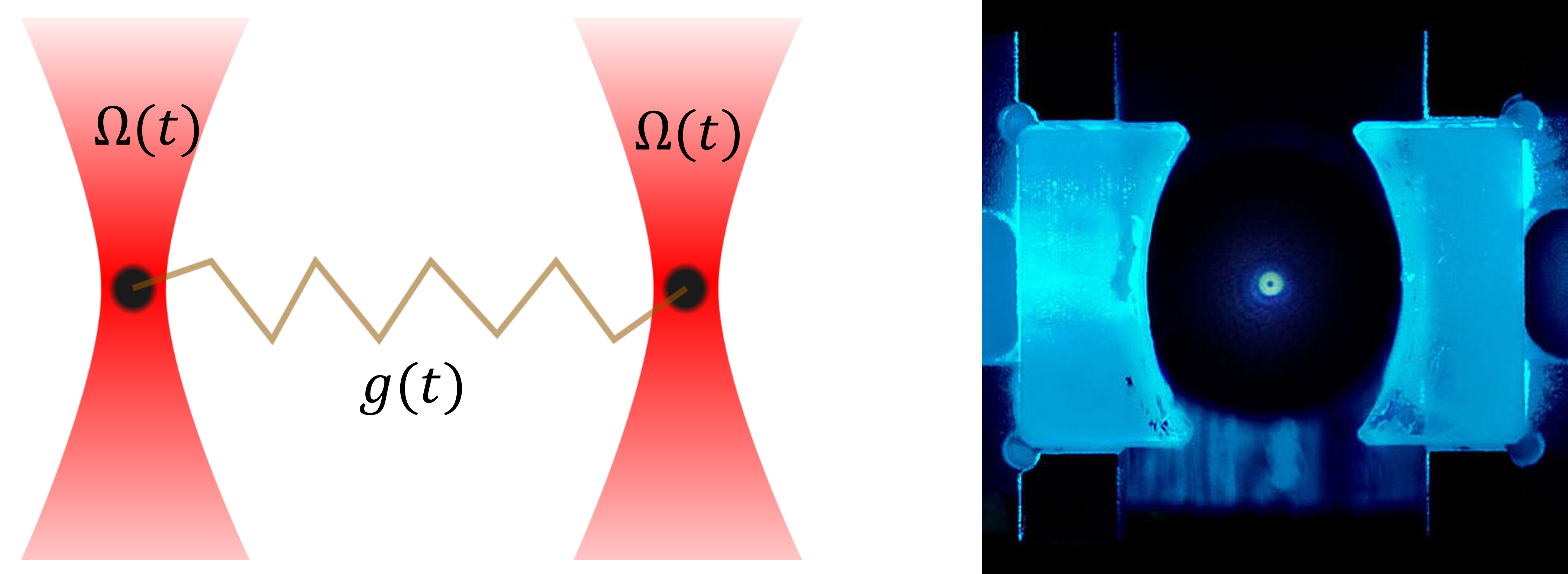


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Abstract

The quantum entanglement of relatively large-massed particles by optomechanics methods has become viable in recent years and garnered interest in various practical and theoretical fields. The Coulombic coupling of two optically levitating charged particles is a potential realization of this system. We focus on building a theoretical framework for such a system by utilizing continuous optical homodyne measurement control and parametric modulation of the particles — a novel and original approach in the field.



Theory – The system

- The system is comprised of two Coulombic interacting harmonically trapped particles.
- The active measurement and control of the system will give rise to the stochastic master equation

$$d\rho_c = -i[H, \rho]dt - \sum_{k=1,2} (\mathcal{D}_{cl}^k[\rho_c] + \mathcal{D}_{pr}^k[\rho_c])dt + \sum_{k=1,2} (\mathcal{H}_{cl}^k[\rho_c])dW_k.$$

- The quadratic master equation ensures that an initial Gaussian state will remain as such. We can therefore focus on the first and second quadrature moments

$$dX^c = A(t)X^c dt + B(t)udt + \Sigma^c C^T(t)W^{-1}dw$$

$$\frac{d\Sigma^c}{dt} = A(t)\Sigma^c + \Sigma^c A(t)^T + V(t) - \Sigma^c C^T(t)W^{-1}C(t)\Sigma^c,$$

where

$$A = \begin{pmatrix} 0 & \Omega(t) & 0 & 0 \\ -\Omega(t) - 2g(t) & -\gamma & 2g(t) & 0 \\ 0 & 0 & 0 & \Omega(t) \\ 2g(t) & 0 & -\Omega(t) - 2g(t) & -\gamma \end{pmatrix}.$$

Theory – Parametric resonance

- One can diagonalize the system into its two eigenfrequencies

$$\Omega_+ = \Omega$$

$$\Omega_- = \sqrt{\Omega^2 + 4g\Omega}.$$

- The system can be parametrically driven by modulating the following parameters:

1. The trap's frequency,

$$\Omega(t) = \Omega_0 + 2\Omega_1 \cos(\Omega_{trap}t).$$

2. The Coulombic coupling strength

$$g(t) = -|g_0| + 2g_1 \cos(\Omega_{coupling}t).$$

- Different modulation values can increase/decrease the entanglement of the two particles.

- The entanglement is characterized by the logarithmic negativity of the system.

Theory – Floquet Theorem

- The presence of periodic time modulation leads to the emergence of exponentially diverging\decaying Floquet modes, governed by the linear equation

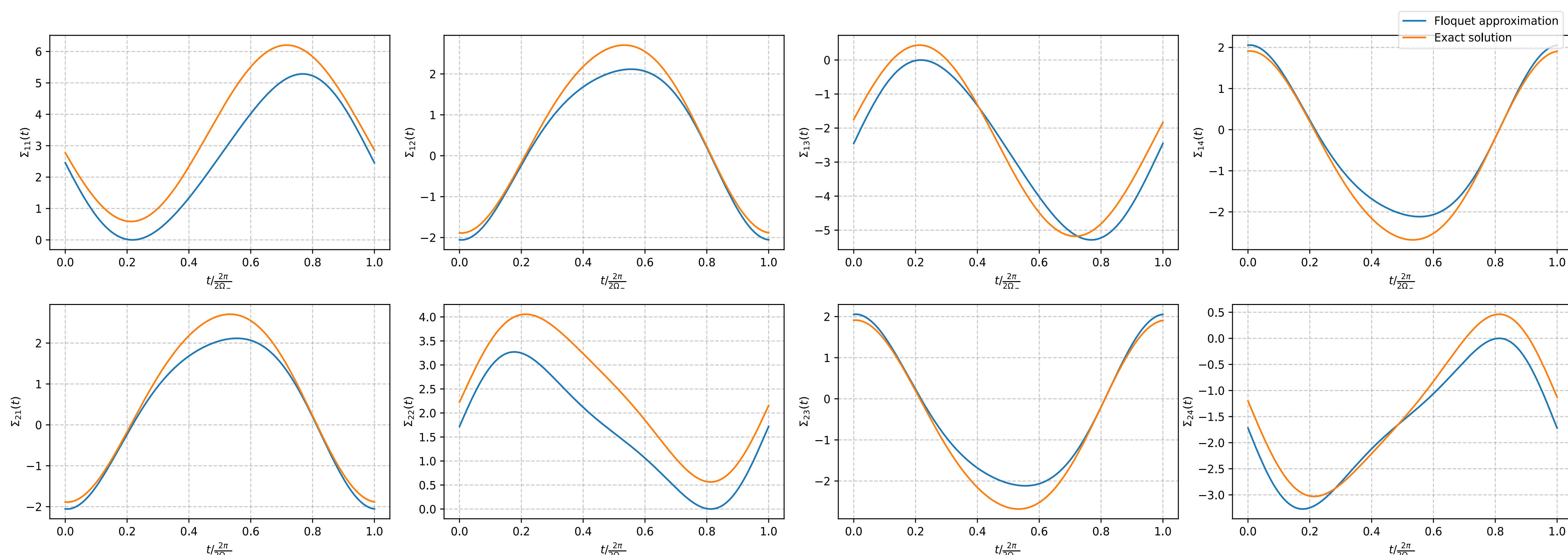
$$\frac{d\Sigma^F}{dt} = A(t)\Sigma^F + \Sigma^F A(t)^T \quad \Sigma_F(t) = e^{\lambda t} \widetilde{\Sigma}_F(t),$$

- The system tends to follow the strongest Floquet mode, while the nonlinearity cancels the divergence.

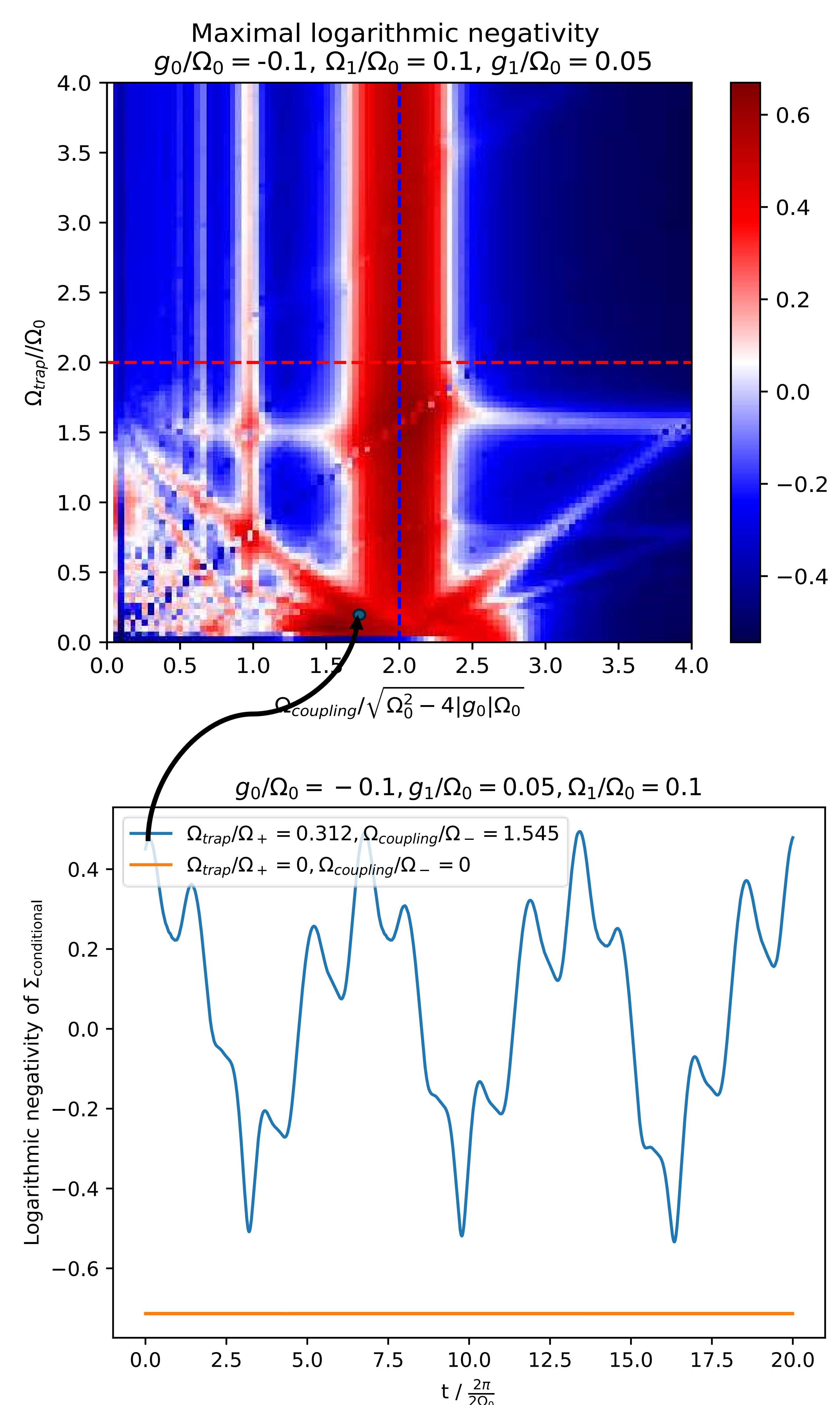
$$\Sigma^c(t) \propto \widetilde{\Sigma}_F(t)$$

- The proportionality constant can be approximately calculated by analytics.

Analytical vs numerical results



Numerical results



References

- [1] Magrini, L., Rosenzweig, P., Bach, C. et al. Real-time optimal quantum control of mechanical motion at room temperature. Nature 595, 373–377 (2021).
- [2] K. Winkler, A. V. Zasedatelev, B. A. Stickler, U. Delic, A. Deutschmann-Olek, and M. Aspelmeyer, Steady-state entanglement of interacting masses in free space through optimal feedback control, arXiv (2024).
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- [3] Magrini, Lorenzo, and Yuriy Coroli. Cooling of a Levitated Nanoparticle to the Motional Quantum Ground State." University of Vienna, 31 Jan. 2020, <https://quantum.univie.ac.at/news/details-news/news/cooling-of-a-levitated-nanoparticle-to-the-motional-quantum-ground-state-1/>.