

Concepts of condensed matter physics - Exercise #4

Spring 2018

Due date: 07/06/2018

1. In this question you will re-derive the BCS theory studied in class and use it to calculate a few properties of superconductors. Our starting point is the Hamiltonian of electrons interacting via an attractive point contact interaction ($g > 0$):

$$\hat{H} = \int d^3x \left[\sum_{s=\uparrow,\downarrow} c_s^\dagger(x) \left(-\frac{\nabla^2}{2m} - \mu \right) c_s(x) - g c_\uparrow^\dagger(x) c_\downarrow^\dagger(x) c_\downarrow(x) c_\uparrow(x) \right].$$

- Write the Hamiltonian in momentum space, and then transform it to a quadratic form by assuming the order parameter $\Delta = \frac{g}{\Omega} \sum_k c_{-k\downarrow} c_{k\uparrow}$ is weakly fluctuating (i.e., by performing mean field). Here Ω is the system's volume.
 - Diagonalize the quadratic Hamiltonian and find the spectrum of excitations.
 - What is the ground state wavefunction? What is the ground state energy? Show that taking $\Delta = 0$ we recover the known non-interacting ground state energy.
 - Using the ground state wavefunction, write a self-consistent equation ("the BSC gap equation") for Δ . Solve this equation for small values of g .
 - Extend the gap equation to finite temperatures by promoting the average with respect to the ground state to a thermal average.
 - Find the critical temperature T_c above which superconductivity is destroyed. What is the value of Δ slightly below the transition?
2. In this question you will find the spectrum of the above BCS theory in the presence of spin-orbit and Zeeman coupling. In **one-dimension**, the Hamiltonian is given by:

$$\hat{H} = \int dx \left[\sum_{s,s'} c_s^\dagger(x) \left(-\frac{\partial_x^2}{2m} + iu\sigma_z^{ss'} \partial_x + B\sigma_x^{ss'} - \mu \right) c_{s'}(x) - g c_\uparrow^\dagger(x) c_\downarrow^\dagger(x) c_\downarrow(x) c_\uparrow(x) \right]$$

- a. First, neglecting g , diagonalize the quadratic Hamiltonian by going to momentum space. Draw the spectrum (qualitatively) – how does the spin-orbit and Zeeman terms alter the parabolic spectrum of free electrons

$$(E = \frac{k^2}{2m} - \mu).$$

- b. Introducing finite g and performing a mean field approximation, write a quadratic Hamiltonian of the form:

$$H = E_0 + \sum_k \vec{\Psi}_k^\dagger h_{BDG}(k) \vec{\Psi}_k,$$

$$\text{With } \vec{\Psi}_k = \begin{pmatrix} c_\uparrow(k) \\ c_\downarrow(k) \\ c_\uparrow^\dagger(-k) \\ -c_\downarrow^\dagger(-k) \end{pmatrix} \text{ and } h_{BDG}(k) = \begin{pmatrix} A(k) & D(k) \\ D^\dagger(k) & -\sigma_y A^*(-k) \sigma_y \end{pmatrix}.$$

Find the 2×2 matrices $A(k)$ and $D(k)$.

- c. Diagonalize h_{BDG} and find the spectrum of excitations.
d. Show that by changing the ratio Δ / B , we reach a point in which the gap to excitations closes. Draw the spectrum at this point.

3. Superconductivity on the surface: In this question you will find that above H_{c2} there is a range of fields for which superconductivity can survive on the surface. Consult “Introduction to superconductivity”, by M. Tinkham, page 135.

- a. Start from the Ginzburg-Landau theory of a superconductor and neglect non-quadratic orders near the critical point. Write down the corresponding equations of motion, and using an analogy to the Schrodinger equation, find the critical field H_{c2} , above which superconductivity cannot nucleate in the interior of the sample. Write the result in terms of ϕ_0 and ξ . Can you explain the result qualitatively?
b. Consider the same physical setting with an edge at $x = 0$ (such that for $x > 0$ there is an insulator). Show that the boundary conditions take the form

$\left(\frac{\nabla}{i} - \frac{2\pi A}{\phi_0}\right)\psi\Big|_n = 0$. Show that one can automatically satisfy this boundary condition by considering an auxiliary potential, containing a mirror image of the original potential in the insulating region. Does this affect the solution from part (a) well inside the superconductor (i.e., for $|x| \gg \xi$)?

- c. Argue, using the auxiliary potential, that very close to the surface one can find a solution with lower energy, making the critical field higher near the surface.